

1 **Estimation of seismic centroid moment tensor using ocean bottom pressure gauges**
2 **as seismometers**

3

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10 **Key Points:**

- 11 • We estimated the CMT of offshore $M \sim 7$ earthquakes using onshore seismometers
12 and offshore pressure gauges
- 13 • The horizontal location of the centroid is well constrained by using offshore pressure
14 gauges as seismometers
- 15 • Observed pressure-change waveforms show the theoretical predicted relationship
16 between pressure and vertical acceleration

17

18 **Abstract**

19 We examined the dynamic pressure change at the seafloor to estimate the centroid moment
20 tensor solutions of the largest and second largest foreshocks (Mw 7.2 and 6.5) of the 2011
21 Tohoku-Oki earthquake. Combination of onshore broadband seismograms and high-
22 frequency (~20–200 s) seafloor pressure records provided the resolution of the horizontal
23 locations of the centroids, consistent with the results of tsunami inversion using the long-
24 period (> ~10 min) seafloor pressure records although the depth was not constrained well,
25 whereas the source locations were poorly constrained by the onshore seismic data alone. Also,
26 the waveforms synthesized from the estimated CMT solution demonstrated the validity of the
27 theoretical relationship between pressure change and vertical acceleration at the seafloor. The
28 results of this study suggest that offshore pressure records can be utilized as offshore
29 seismograms, which would be greatly useful for revealing the source process of offshore
30 earthquakes.

31

32 **1. Introduction**

33 When earthquakes occur offshore, associated tsunami are observed by ocean-bottom
34 pressure gauges (OBPGs). For example, the DART (Deep-ocean Assessment and Reporting
35 of Tsunamis) systems developed by NOAA (National Oceanic and Atmospheric
36 Administration) were founded on a network of widely distributed OBPGs to monitor tsunami
37 far offshore (e.g., González et al., 2005). OBPGs are considered one of the most reliable
38 sensors for investigating tsunami propagation and source models because they are free from
39 the strong site effects usually observed near coasts. Near-field (less than ~ 100 km from the
40 epicenters) OBPG tsunami records enable us to obtain good spatial resolution of the source
41 models of offshore moderate ($M \sim 7$) earthquakes (e.g., Saito et al., 2010; Kubota et al.,
42 2015; 2017). Lack of resolution for source models will be a barrier for studying the detailed
43 source processes of offshore earthquakes (e.g., Heiderzadeh et al., 2017a).

44 Near-field pressure records obtained ~ 20 km from two local earthquakes (Mw 7.2 and
45 Mw 6.5, National Research Institute for Earth Science and Disaster Resilience [NIED],
46 2011a; 2011b) (the station and earthquake locations are in Figure 1) are shown in Figure 2.
47 By applying a low-pass filter (>400 s) to the original records (gray), we obtained clear
48 tsunami signals (blue). The maximum tsunami height was ~ 10 cm (Figure 2b) and ~ 1 cm
49 (Figure 2e) for the Mw 7.2 and 6.5 earthquakes, respectively. Some studies have estimated
50 earthquake fault models by analyzing such tsunami signals in the OBPGs (e.g., Gusman et al.,
51 2013; Heidarzadeh et al., 2017b; Kubota et al. 2017).

52 In addition to tsunami, OBPGs can observe other signals associated with earthquakes.
53 When OBPGs are installed inside the focal area, permanent seafloor vertical deformations are
54 observed as the difference between the average pressure levels before and after the
55 earthquake. These are often used to estimate coseismic fault models (e.g., Ito et al., 2011;
56 2013; Iinuma et al., 2012; Ohta et al., 2012; Tsushima et al., 2012; Wallace et al., 2016). The
57 pressure changes associated with tsunami and permanent deformations are interpreted as the
58 change in the loading due to the water column over the OBPG based on the hydrostatic
59 assumption. OBPGs also observe dynamic pressure changes associated with seismic waves
60 (e.g., Filloux, 1982; Bolshakova et al., 2011; Matsumoto et al., 2012; Saito & Tsushima,
61 2016), which are caused by seafloor seismic motions and ocean acoustic waves. In Figure 2,
62 high-frequency pressure changes are evident from large amplitudes, especially after applying

63 a bandpass filter with a passband of 0.01–0.05 Hz (red lines). These components are usually
 64 removed before tsunami waveform analyses (e.g., Gusman et al., 2013; Inazu & Saito, 2014;
 65 Kubota et al., 2015; 2017; Heidarzadeh et al., 2016), because they are irrelevant to the sea
 66 surface displacement due to tsunami.

67 The dynamic pressure change associated with seismic motion has been previously
 68 studied (e.g., Filloux, 1982; Nosov & Kolesov, 2007; Bolshakova et al., 2011; Matsumoto et
 69 al., 2012; Saito, 2013). The pressure changes can be interpreted based on two different
 70 relationships according to their frequency range, defined by the fundamental acoustic
 71 resonant frequency $f_0 = c_0/4h_0$ (h_0 is the water depth and c_0 is the velocity of the ocean
 72 acoustic wave). When the frequency of the seafloor motion is sufficiently low compared to
 73 the fundamental acoustic resonant frequency f_0 ($f < f_0$), the seafloor pressure change can be
 74 approximated as:

$$75 \quad p = \rho_0 h_0 a_z, \quad (1)$$

76 and when the frequency is high ($f > f_0$) as:

$$77 \quad p = \rho_0 c_0 v_z, \quad (2)$$

78 where ρ_0 is seawater density and a_z and v_z are the vertical acceleration and velocity of the
 79 seafloor motion (hereafter, pressure–acceleration relationship and pressure–velocity
 80 relationship), respectively (e.g., Bolshakova et al., 2011; Matsumoto et al., 2012). Numerical
 81 simulation is useful for investigating these relationships (e.g., Maeda et al., 2013; Kozdon &
 82 Dunham, 2014; Saito & Tsushima, 2016; Saito, 2017). Saito and Tsushima (2016) tried to
 83 reproduce the dynamic pressure change associated with the 2011 Tohoku-Oki earthquake by
 84 numerical simulation, assuming a uniform slip fault model. The simple model roughly
 85 reproduced the dynamic pressure changes, but not completely. Using the seismic equations
 86 considering a compressible sea and elastic crust, Saito (2017) numerically simulated the
 87 vertical acceleration and pressure at the seafloor, and found that the pressure–acceleration
 88 relationship (equation (1)) works well, whereas the pressure–velocity relationship (equation
 89 (2)) works only for the first motion of the pressure change.

90 Although many previous studies have investigated the pressure–acceleration
 91 relationship (equation (1)) based on theoretical studies or numerical simulations, there are
 92 few studies based on real observations. Matsumoto et al. (2012) showed the pressure–

93 acceleration relationship works reasonably well in the frequency domain by comparing
94 records of the Tohoku-Oki earthquake observed ~400 km away from the focal area. Nosov
95 and Kolesov (2007) also investigated pressure records of the 2003 Tokachi-Oki earthquake
96 (Mw 8.0), but only in the frequency domain. Those studies did not compare records in the
97 time domain. Hence, it is not confirmed whether the phases are in agreement or have some
98 shift in the time domain. If the pressure–acceleration relationship works well in the time
99 domain, we could estimate various earthquake parameters by applying seismological analyses
100 to the pressure data.

101 The purpose of this study is to clarify whether the pressure–acceleration relationship
102 holds in observed records through a centroid moment tensor (CMT) analysis of offshore
103 earthquakes. Moreover, we demonstrate that the use of offshore OBPB records improves the
104 centroid horizontal locations of the offshore earthquakes.

105

106 **2. Ocean Bottom Pressure Data**

107 We used pressure data associated with the Mw 7.2 and 6.5 interplate earthquakes off
108 northeastern Japan on 9 March, 2011, which were the largest and second largest events
109 preceding the 2011 Tohoku-Oki earthquake (Ohta et al., 2012; Gusman et al., 2013; Kubota
110 et al., 2017) (hereafter foreshock #1, foreshock #2, and the mainshock, respectively). Since
111 the magnitudes of the two foreshocks were large enough to show good signal-to-noise ratios,
112 the source processes are relatively simple compared to the mainshock, and we can expect to
113 obtain a reasonable CMT solution with a point-source assumption. Also, the rupture areas of
114 these earthquakes have been estimated from tsunami data by Kubota et al. (2017) (colored
115 rectangles in Figure 1), and those estimates can be used as a reference for validating the
116 centroid location.

117 We used 1 Hz sampled pressure data. Seven offline autonomous OBPGs with pop-up
118 recovery were deployed within ~70 km of the epicenters (Figure 1). Details are given in Hino
119 et al. (2014). Real-time cabled OBPGs of the Earthquake Research Institute of the University
120 of Tokyo (Kanazawa & Hasegawa, 1997), TM1 and TM2, were also in operation more than
121 100 km away from the epicenters (Figure 1).

122 Figure 3 shows the original and bandpass-filtered pressure records. We applied the
 123 Butterworth-type bandpass filter in both the forward and reverse directions. The passband
 124 was determined according to the fundamental resonant frequency f_0 and the frequency range
 125 used in the F-net Moment Tensor (MT) analysis by the National Research Institute for Earth
 126 Science and Disaster Resilience (NIED, 2011a; 2011b). Supposing the acoustic wave velocity
 127 $c_0 = 1.5$ km/s and the water depth $h_0 = 1.5$ km (average depth of the focal area), the
 128 fundamental resonant frequency ($f_0 = c_0/4h_0$) is ~ 0.25 Hz (4 s). If the sea depth is 3.2 km
 129 (corresponding to the depth of the deepest OBPG, GJT3), f_0 is ~ 0.12 Hz (8.5 s). Since the
 130 pressure–acceleration relationship holds when the dominant frequency is lower than the
 131 acoustic resonant frequency ($f < f_0$), the high-frequency component ($f > \sim 0.1$ Hz) should be
 132 suppressed. As for the low-frequency cutoff, tsunami components ($T > \sim 400$ s, $f < \sim 0.0025$
 133 Hz) should be reduced. The F-net MT analysis adopted a passband of 0.005–0.02 Hz (50–200
 134 s) for foreshock #1, and 0.01–0.05 Hz (20–100 s) for foreshock #2 (NIED, 2011a; 2011b).
 135 Considering the factors above, we took the passband to be 0.005–0.02 Hz for foreshock #1,
 136 and 0.01–0.05 Hz for foreshock #2.

137

138 **3. Methodology and Results: Centroid Moment Tensor Inversion using Dynamic** 139 **Pressure Records**

140 We estimated the centroid moment tensors, centroid times, and centroid locations of
 141 the two earthquakes by analyzing onshore seismic and offshore dynamic pressure data. The
 142 procedure for the CMT inversion followed the grid-search approach of Ito et al. (2006),
 143 which uses five independent basis MT components (Kikuchi & Kanamori, 1991) (details of
 144 the calculation of Green's functions and the CMT inversion are given in Text S1). In the
 145 analysis, the seafloor vertical acceleration is calculated using a conventional elasto-dynamic
 146 equation using the discrete wavenumber method (e.g., Saikia, 1994) with the 1-D subsurface
 147 structure model of Kubo et al. (2002), and is converted to the dynamic pressure change using
 148 the pressure–acceleration relationship (equation (1)), where the water density ρ_0 is set as 1.03
 149 g/cm³. The sea depths h_0 of the OBPGs are summarized in Table S1. The same bandpass
 150 filter used for the observation is applied to the calculated waveforms. As a measure of
 151 waveform reproducibility, we used variance reduction (VR):

$$152 \quad VR = \left(1 - \frac{\sum_{k=1}^N (w_k d_k^{obs} - w_k d_k^{calc})^2}{\sum_{k=1}^N (w_k d_k^{obs})^2} \right) \times 100 \text{ (\%)} \quad (3)$$

153 where d_k^{obs} and d_k^{calc} are the k -th data of observed and calculated waveforms, respectively,
 154 and N denotes the number of data used for inversion. Note that the onshore seismometers and
 155 offshore OBPBs have different dimensions. To reduce the bias caused by the difference in
 156 the inversion analysis, we introduced the weight value, w_k , as the inverse of the maximum
 157 amplitude of each waveform:

$$158 \quad w_k = \frac{1}{\max(d_i(t))}, \quad (4)$$

159 where $d_i(t)$ is the time series of the i -th station including the k -th datum. With respect to the
 160 grid search, we sought the horizontal and vertical locations of the centroid at 0.1° intervals in
 161 the horizontal and 2 km in the vertical. The interval of the temporal grids is 1 s.

162 In the analysis, we used two types of datasets, as follows. One of the datasets consists
 163 only of the onshore seismograms (dataset 1). Three (radial, transverse, and vertical) velocity
 164 components were obtained by the F-net stations (e.g., Okada et al., 2004). We analyzed the
 165 same datasets used in the F-net MT solution: the F-net stations NOP, WJM, and WTR for
 166 foreshock #1 (NIED, 2011a, red triangles in Figure 1) and IMG, KZK, and SGN for
 167 foreshock #2 (NIED, 2011b, blue triangles). We prepared another dataset (dataset 2) by
 168 adding the pressure data obtained at GJT3, which is located on the offshore side of the
 169 epicenters, to dataset 1, in order to improve the station coverage.

170 The estimated CMT solutions of foreshock #1 are shown in Figure 4. Using only the
 171 onshore seismograms (dataset 1), the centroid of foreshock #1 was estimated at 143.2°E ,
 172 38.4°E , and 30 km with a VR of 66.1% (gray CMT solution in Figure 4a). The strike, dip,
 173 and rake were 171.2° , 21.2° , and 52.0° , respectively, and the seismic moment M_0 was $5.5 \times$
 174 10^{19} Nm (Mw 7.1). This solution was close to the F-net solution in location and mechanism
 175 (143.2798°E , 38.3285°E , 23 km, and 8.9×10^{19} Nm; black CMT solution in Figure 4a)
 176 (NIED, 2011a). The centroid was not located inside the rupture area estimated by the tsunami
 177 waveform analysis (Kubota et al., 2017) (red rectangle in Figure 4b).

178 On the other hand, using dataset 2 (pressure data from GJT3 included), the centroid
 179 was estimated at 142.9°E , 38.5°E , and 30 km (red CMT solution in Figure 4b) with a VR of

180 61.9%; the strike, dip, and rake were 164.1° , 23.2° , and 44.1° , respectively, and $M_0 = 5.1 \times$
181 10^{19} Nm (M_w 7.1). The centroid was located almost at the center of the rupture area obtained
182 from tsunami waveform inversion (Kubota et al., 2017).

183 We conducted forward simulations of OBPG waveforms not used for the inversion, to
184 examine the agreement with and observation. Both CMT solutions estimated from dataset 1
185 and dataset 2 reproduced the seismograms of the F-net stations nicely (Figure 4c), suggesting
186 that the difference in the centroid horizontal locations cannot be resolved using only the
187 seismograms of onshore stations. In contrast, we recognize the difference in the pressure
188 waveforms recorded at offshore stations. The waveforms obtained from dataset 2 (red lines in
189 Figure 4d) reproduce the observations (black) better than those from dataset 1 (gray). Note
190 that we used only GJT3 in the inversion analysis, but we also see this improvement in other
191 pressure records. For example, if we evaluate the VR of P09 (VR_{P09}) using the same time
192 window used for inversion (white background area in Figure 4d), we obtained 2.5% for the
193 CMT solution from dataset 1 and 70.6% for the solution of dataset 2.

194 We obtained similar results for foreshock #2. The centroid location of foreshock #2
195 obtained from dataset 1 was $143.2^\circ E$, $38.2^\circ E$, and 26 km with a VR of 80.1% (gray CMT
196 solution in Figure 5a), and the strike, dip, and rake were 194.0° , 19.1° , and 77.3° ,
197 respectively ($M_0 = 4.2 \times 10^{18}$ Nm, M_w 6.4). These are similar to the F-net solution (black) of
198 $143.0448^\circ E$, $38.1722^\circ E$, and 20 km ($M_0 = 5.51 \times 10^{18}$ Nm) (NIED, 2011b). When GJT3 was
199 included (dataset 2), the centroid was at $142.9^\circ E$, $38.3^\circ E$, and 32 km with a VR of 68.8%,
200 and the strike, dip and rake were 144.1° , 39.2° , and 21.1° , respectively ($M_0 = 4.2 \times 10^{19}$ Nm,
201 M_w 6.4) (Figure 5b). We found that the centroid estimated from dataset 2 was closer to the
202 center of the fault model than that from dataset 1. Also, the CMT solution estimated using
203 dataset 2 reproduced the OBPG records better: the values of VR_{P09} were -13.6% (dataset 1)
204 and 48.0% (dataset 2).

205 To evaluate the resolution and accuracy of the centroid horizontal location, we
206 calculated the area where the VR exceeds 90% of the best-fit VR in each result. The high-VR
207 area ($>90\%$) is surrounded by gray lines in Figures 4a, 4b, 5a, and 5b. For both foreshocks #1
208 and #2, the high-VR area extended in the ENE–WSW direction when pressure records from
209 GJT3 were not used for the inversion (dataset 1) by ~ 100 km and ~ 50 km, respectively
210 (Figures 4a and 5a), suggesting the horizontal location of the centroid is not well constrained.

211 On the other hand, the high-VR area obtained from dataset 2 was much smaller and the EW
212 extent of the high-VR area became by half (by ~50 km for foreshock #1 and ~25 km for
213 foreshock #2), and mostly confined within the spatial extent of the finite fault models derived
214 by the tsunami inversion (Figures 4b and 5b). Since the station coverage was improved by
215 adding the dynamic pressure obtained at GJT3, located on the offshore side of the focal area
216 (Figure 1), the resolution of the centroid horizontal location was significantly improved.

217 The obtained centroid depths (~30 km) are systematically deeper than the plate
218 boundary depth obtained by the seismic survey (Ito et al., 2005). To discuss the depth
219 resolution, we examined the vertical VR distribution at the horizontal point where the best-fit
220 CMT is obtained (Figures S1 and S2). When the centroid depth is less than ~10 km, the VR
221 is smaller than 90% of the maximum VR, suggesting that the centroid depth is deeper than 10
222 km. However, the vertical range where the VR exceeds 90% of the best value (gray lines in
223 Figures S1 and S2) is widely distributed, suggesting the centroid depths are not well
224 constrained. We also point out that that the centroid depths and the centroid time delay have a
225 trade-off relationship.

226

227 **4. Discussion**

228 **4.1. Checking the validity of the pressure–acceleration relationship**

229 Past studies have found that the seismograms and pressure records in the Fourier
230 amplitude are consistent with the theoretical relationship given by equation (1) (e.g.,
231 Matsumoto et al., 2012). However, equation (1) suggests not only an agreement in the
232 Fourier amplitude but also in the time series of the accelerograms and pressure records. It is
233 important to confirm the agreement in the waveforms by observations. Our results confirmed
234 equation (1) by indirect comparison of the waveforms: the observed pressure records and
235 theoretically calculated accelerograms. In Figures 4d and 5d, we compare the observed
236 pressure records and pressure changes converted from the calculated accelerograms based on
237 equation (1). The theoretically predicted pressure changes (red and blue lines) agree well
238 with the observations (gray lines) at stations P02, P06, and P09, although those data were not
239 used in the inversion analysis. Hence, the agreement found in those waveforms strongly
240 supports the validity of equation (1) in time domain. Disagreement at station P08 may be

241 caused by insufficient modeling using the point-source approximation near the finite-fault
242 source.

243

244 **4.2. Importance of offshore dynamic pressure change for source estimation**

245 Our results show that OBPGs are very useful in constraining centroid horizontal
246 locations. This is mainly because of the improvement in station coverage achieved by using
247 the pressure data as offshore seismograms. A high sampling rate (1 Hz) also contributes to
248 obtaining the CMT solution. Filloux (1982) has already suggested the applicability of
249 seafloor pressure data as seismometers, but the sampling rate in his analysis was very low (28
250 s). If the sampling rate is low, aliasing due to the higher-frequency ocean acoustic wave (with
251 a dominant period of ~5–10 s) may prevent us from obtaining high-quality bandpass filtered
252 waveforms. The use of 1-Hz-sampling pressure data enabled us to use high-quality
253 seismograms for CMT analysis. The dynamic pressure records, free from amplitude
254 saturation, can be treated as on-scale near-field seismic records of offshore large earthquakes,
255 as demonstrated here. As mentioned by Heidarzadeh et al. (2017a), lack of resolution for
256 source models of offshore moderate earthquakes will make it difficult to investigate the
257 detailed source processes. It is expected that broadband seismograms provided by OBPGs
258 will contribute to estimating earthquake source processes in future studies, such as the finite
259 fault model or source duration. The source duration would be useful for identifying tsunami
260 earthquakes. The investigation of more examples would be necessary to confirm the
261 applicability of OBPGs as the broadband seismometer in more detail (e.g., lower limit of the
262 analyzable magnitude range).

263

264 **5. Conclusions**

265 Using the dynamic pressure data observed by OBPGs deployed near the focal area
266 together with onshore seismograms, we estimated CMT solutions for two moderate ($M \sim 7$)
267 offshore interplate earthquakes, and evaluated the robustness of the estimation of the
268 horizontal centroid location. When offshore OBPG data were excluded, the horizontal
269 location of the centroid was not well resolved, and the best-fit centroid was estimated outside
270 of the rupture areas obtained by a previous study. Meanwhile, the centroid locations were

271 well resolved and reasonably constrained inside the rupture area when we used both onshore
272 seismograms and offshore OBPGs, although the depth resolution were not so good in the
273 present study. The extent of the area where relatively high VRs were obtained (i.e.,
274 estimation error) became around half by adding offshore OBPG record. Also, by using the
275 estimated CMT solution, we successfully simulated the actual observed ocean-bottom
276 pressure change records, except for a few stations near the source. The results of our study
277 indicate that the theoretical relationship of equation (1) is valid for actual observations, and
278 we can improve the source estimation of offshore earthquakes using OBPGs based on
279 equation (1).
280

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293 National Research Institute for Earth Science and Disaster Resilience (NIED), Japan
294 (<http://www.fnet.bosai.go.jp/top.php?LANG=en>). The OBPG data used in this study are
295 available in the supplementary datasets S1, S2, and S3 (the description of the datasets is in
296 Text S2).

297

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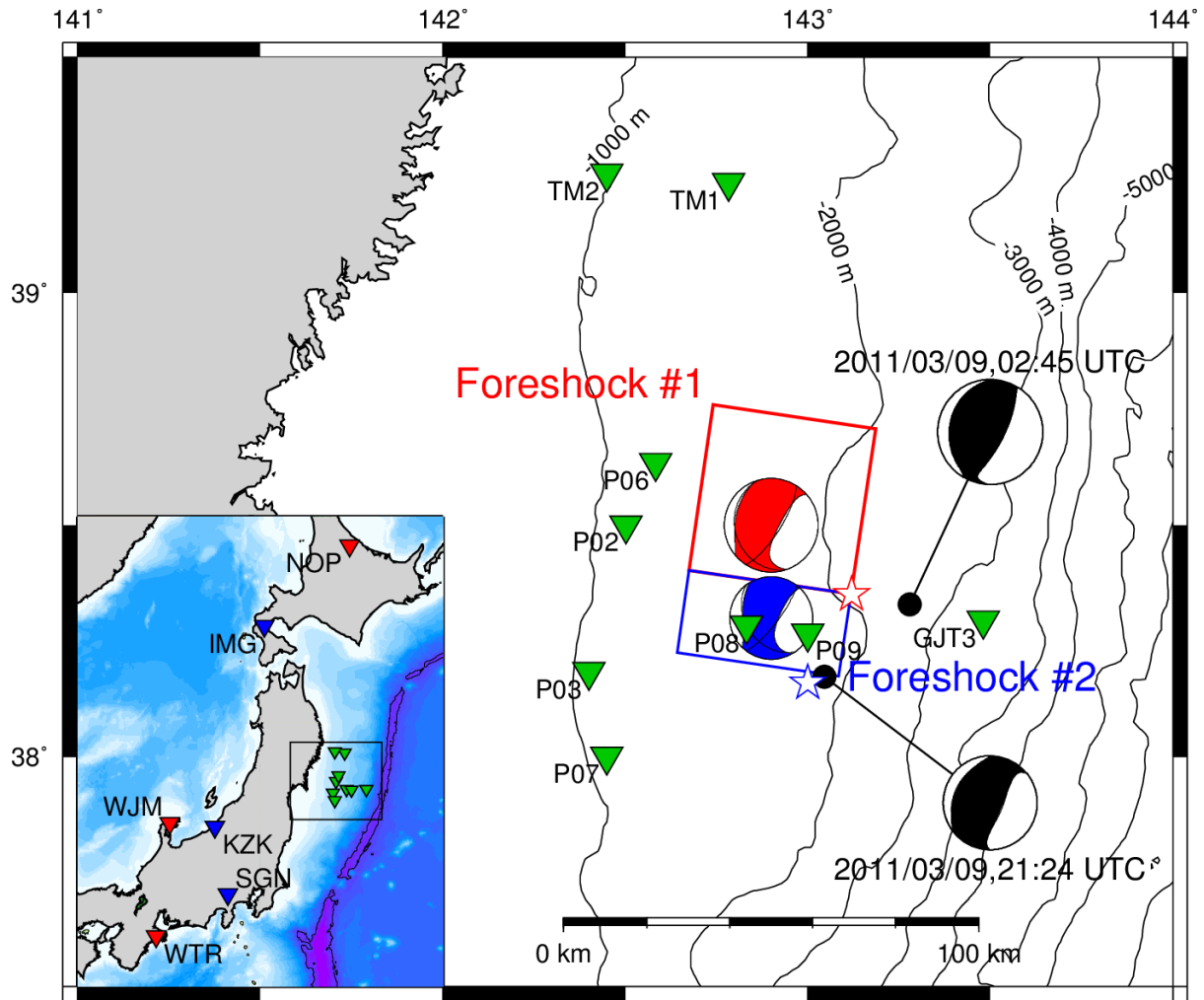
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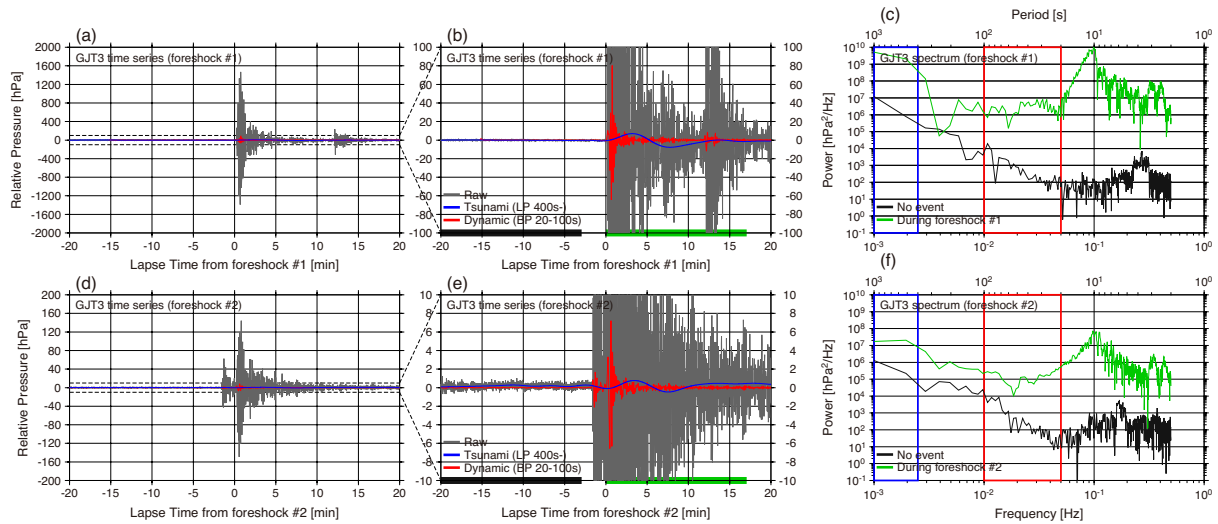
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Figure 1. Location map of this study. Stars are the epicenters determined from ocean bottom seismographs (Suzuki et al., 2012), and rectangles are coseismic rupture areas from near-field tsunami analysis (Kubota et al., 2017) (red: foreshock #1, blue: #2). Inverted triangles denote station locations (green: OBPBs, red and blue: F-net broadband seismometers). F-net routine MT solutions (NIED, 2011a; 2011b) are shown in black, and colored CMT solutions are those obtained jointly using onshore and offshore datasets (also shown in Figures 4b and 5b).



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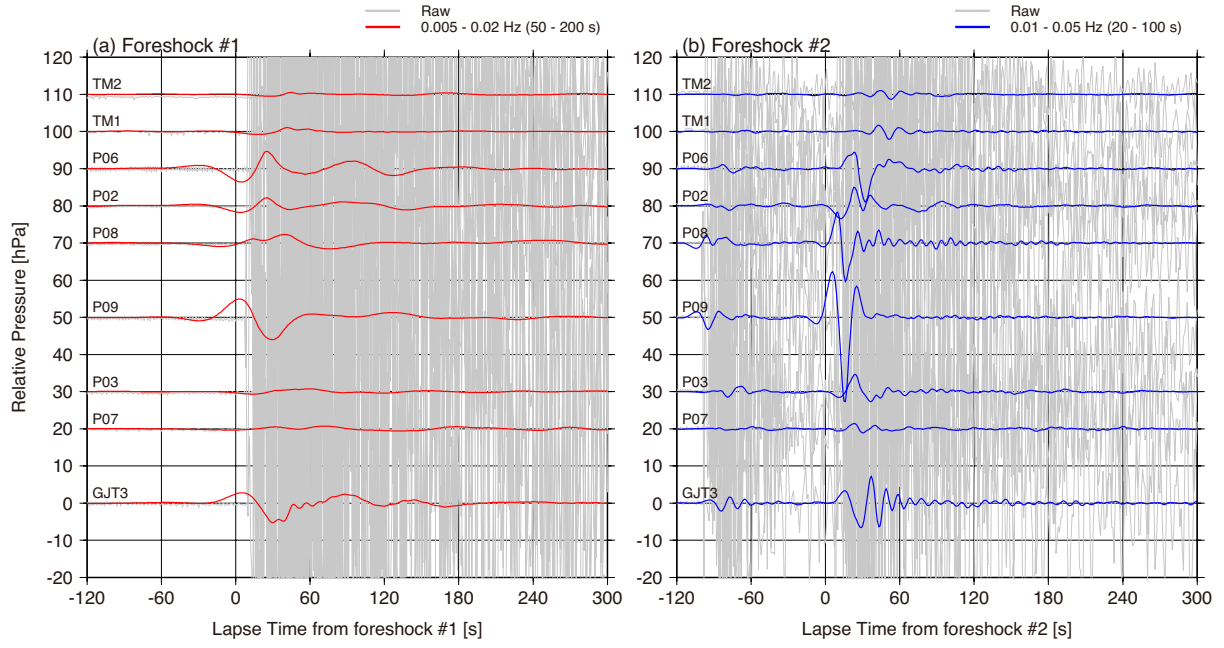
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Figure 2. OBPB records of local earthquakes on 9 March 2011 observed at station GJT3 (shown in Figure 1). (a–b) Time series associated with the Mw 7.2 earthquake. Gray, blue, and red lines are the original, low-pass (>400 s), and bandpass filtered (0.01–0.05 Hz) records, respectively. (c) Power spectra during the Mw 7.2 earthquake (green) and calm period (black), calculated from 1024 s time windows marked by colored bars in Figure 2b. Passbands of the filter in Figure 2b are marked by colored rectangles. (d–f) Time series and power spectra of the Mw 6.5 earthquake.

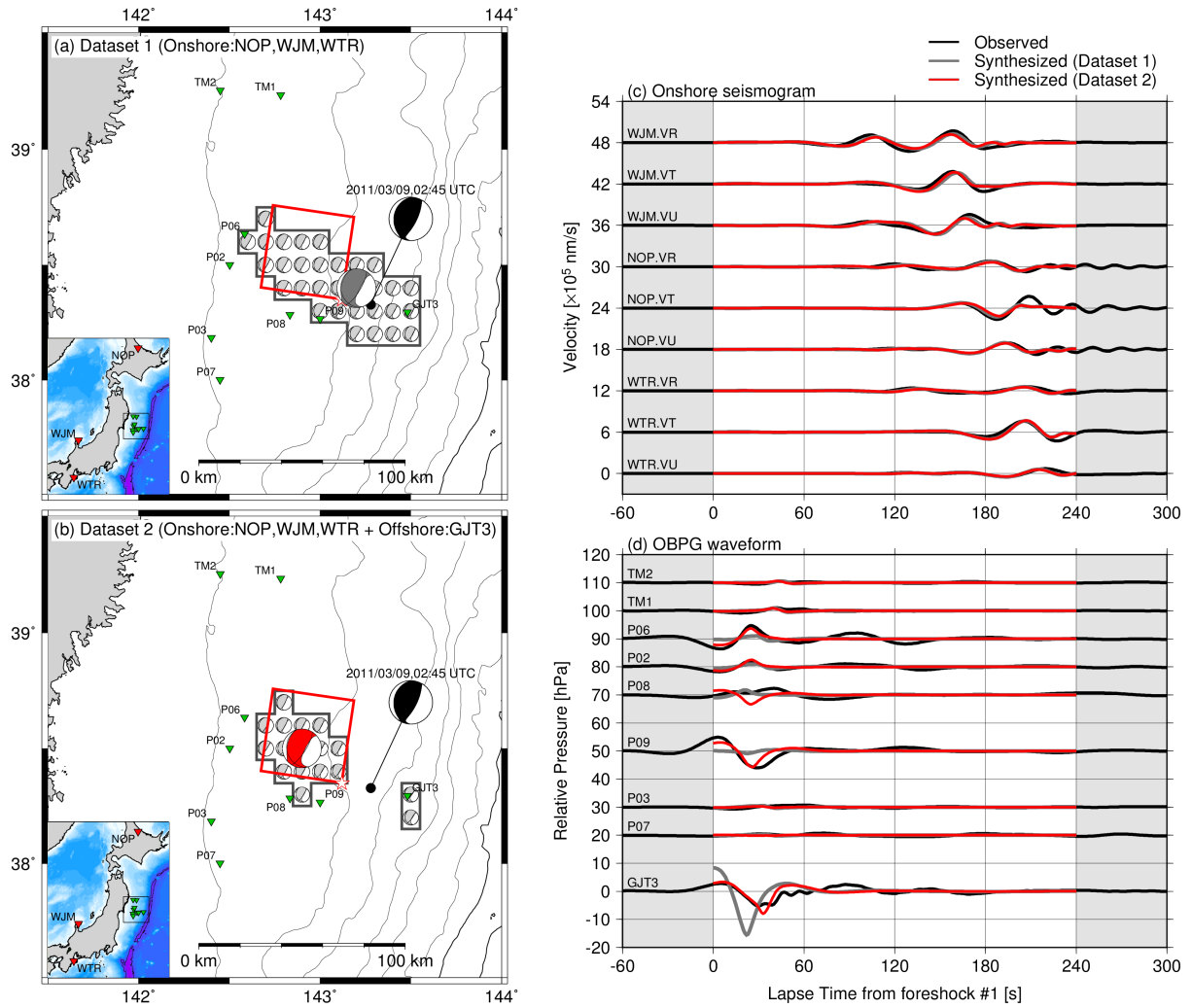


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Figure 3. Comparison of raw pressure waveforms (gray) and bandpass filtered waveforms for (a) foreshocks #1 and (b) #2.



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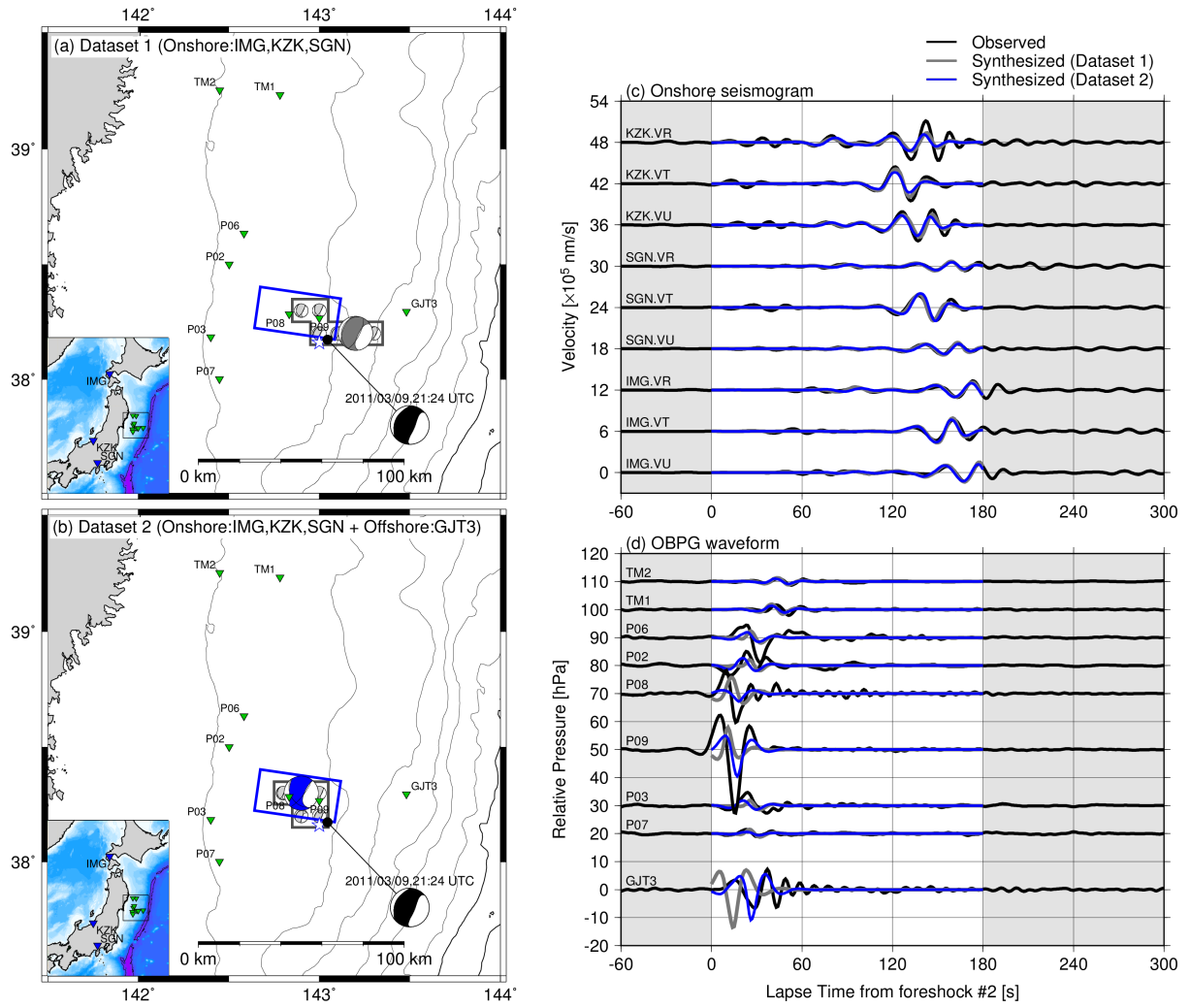
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Figure 4. CMT inversion of foreshock #1. (a) Result obtained from dataset 1, only onshore seismometers (NOP, WJM, and WTR). The best-fit solution is in gray. (b) Result obtained from dataset 2, jointly using onshore seismometers and OBPGs (NOP, WJM, WTR, and GJT3). The best-fit solution is in red. Thick gray lines and small CMTs denote area where the calculated VR exceeds 90% of the best-fit VR. (c) Comparison of onshore seismograms. (d) Comparison of OBPG waveforms, between observed waveforms (black), and synthesized waveforms calculated from the best-fit solution obtained from datasets 1 (gray) and 2 (red), respectively. A time window of 0–240 s (white background area) was used for inversion.



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Figure 5. CMT inversion of foreshock #2. Symbols and colors are the same as in Figure 4.

**Estimation of seismic centroid moment tensor
using ocean bottom pressure gauges as seismometers**

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Introduction

Text S1 describes the method to estimate the centroid moment tensor by waveform inversion. The description for the supplementary datasets (Datasets S1, S2, and S3) is in Text S2. Figure S1 is the vertical distribution of moment tensor solutions to evaluate the resolution of the centroid depth. Table S1 gives the locations of the OBPG stations. The seismic structure used for calculation of the dynamic pressure is shown in Table S2.

Text S1.

In this study, the centroid moment tensors, centroid times, and centroid locations were estimated using the grid-search approach of Ito et al. (2006), under the assumption that the source can be represented by a point source. Here, the detail of the procedure is described.

By assuming the target events is represented by a pure deviatoric moment tensor (MT) without an isotropic component, the observation equation is expressed as:

$$\mathbf{w}^T \mathbf{d} = \mathbf{w}^T \mathbf{G} \mathbf{m}, \quad (\text{S1})$$

where \mathbf{w} is the vector representing the weight of the data, \mathbf{d} is the data vector, \mathbf{G} is the matrix composed of the Green's function, and \mathbf{m} is the model parameter vector consisting of five independent basis MT components (e.g., Kikuchi & Kanamori, 1991). Note that the onshore seismometers and offshore OBPGs have different dimensions. To reduce the bias caused by the difference in the inversion analysis, we introduced the weight value w_k (k -th datum of the weight vector \mathbf{w}):

$$w_k = \frac{1}{\max(d_i(t))}, \quad (\text{S2})$$

where $d_i(t)$ is the time series of the i -th station including the k -th datum. From equation (S1), we obtain the model parameter vector as:

$$\mathbf{m} = [\mathbf{G}^T \mathbf{w} \mathbf{w}^T \mathbf{G}]^{-1} \mathbf{G}^T \mathbf{w} \mathbf{w}^T \mathbf{d}. \quad (\text{S3})$$

We used a time window of 0–240 s from the focal time determined from the ocean bottom seismometers (Suzuki et al., 2012) for foreshock #1, and 0–180 s for foreshock #2, taking their magnitudes and the time windows used in the F-net MT analysis into account.

In the analysis, we calculated the Green's function using the discrete wavenumber frequency method with a 1-D subsurface structure (e.g., Saikia, 1994). Table S2 gives the seismic velocity, attenuation, and density structure used for the calculation, which are the same as those used in the F-net moment tensor calculation, and considered suitable for the 1-D structure of inland Japan (Kubo et al., 2002). Note that we did not assume the effect of the sedimentary layer and the topography for simplicity (i.e., OBPGs are assumed to be located on hard rock on the sea surface). We assumed an impulsive source time function, and the bandpass filter is applied as that used for the observation. Finally, the calculated seafloor vertical acceleration is converted to the dynamic pressure change using the pressure–acceleration relationship $p = \rho_o h_o a_z$ (equation (1)), assuming the water density ρ_o is 1.03 g/cm³, where a_z is the vertical acceleration. The water depth (h_o) of the OBPGs are summarized in Table S1. The same bandpass filters used in the dynamic pressure records are also applied to the Green's function. After we obtained the best fit CMT solution, we forwardly calculated the waveforms which are not used for the inversion analysis to compare with the observation, using the superposition of the Green's functions calculated from five independent basis MT components.

Text S2.

The 1-s sampled raw pressure data used in this study are available in supplementary datasets S1, S2, and S3. This text describes the contents of the datasets.

Dataset S1 is the raw pressure time series for both foreshocks at GJT3, with the time window of -20 min to 20 min from the focal time. This dataset was used to prepare Figure 2. Note that both datasets contain the ocean tide components, although tides were removed in the time series shown in Figure 2 (detail of the tide removal procedure is in Kubota et al. (2017)). The first column denotes the lapse time from the focal time (02:45:16 UTC on 9 March 2011 for foreshock #1 and 21:24:01 UTC for foreshock #2), determined by Suzuki et al. (2012). Times of day (hour, minute, and second in UTC) for both events are also shown.

Datasets S2 and S3 are the raw pressure data for foreshocks #1 and #2, respectively. These datasets were used to prepare Figures 3, 4, and 5. The formats of the time stamps are the same as Dataset S1. The names of the OBPG stations are shown in the first row.

We note that the deployment and retrieval of the OBPGs at GJT3, Po2, Po3, Po6, Po7, Po8, and Po9 were conducted by Tohoku University (Hino et al., 2014; Kubota et al., 2017), and the real-time cabled OBPGs at TM1 and TM2 were operated by Earthquake Research Institute (ERI) of the University of Tokyo (Kanazawa & Hasegawa, 1998). The TM1/TM2 data were resampled to 1 s, although the sampling rate of the original ones was 10 Hz.

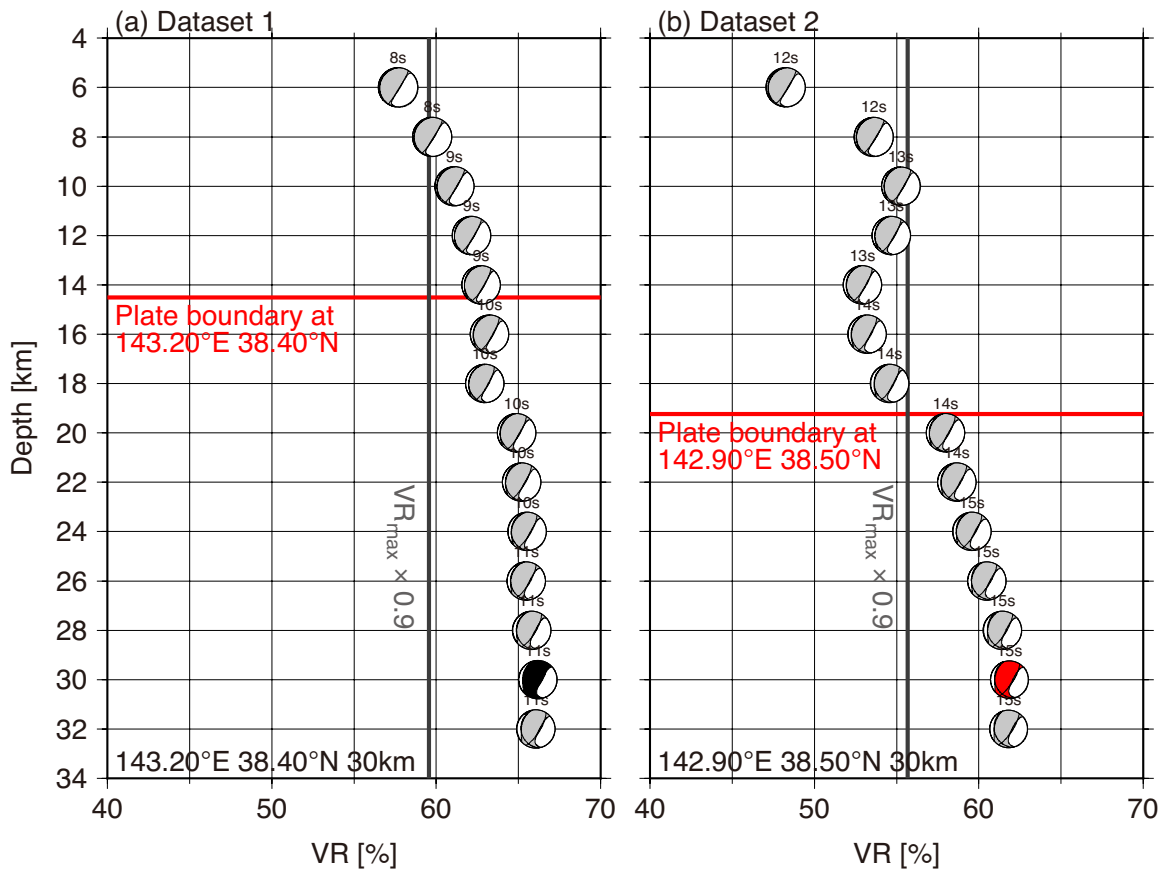


Figure S1. Vertical distribution of VRs and moment tensors at the horizontal location of the best-fit CMT solution for foreshock #1 (Figure 4). (a) Result from dataset consisting of only the onshore seismograms (Figure 4a). (b) Result from dataset consisting of onshore seismograms and offshore pressure data (Figures 4b). The location of the best-fit centroid is shown in the bottom left in each figure. Horizontal and vertical axes denote VR and centroid depth, respectively. Small numbers above each solution are the centroid time delay from the focal time. Red lines denote plate boundary depths obtained from seismic surveys by Ito et al. (2005).

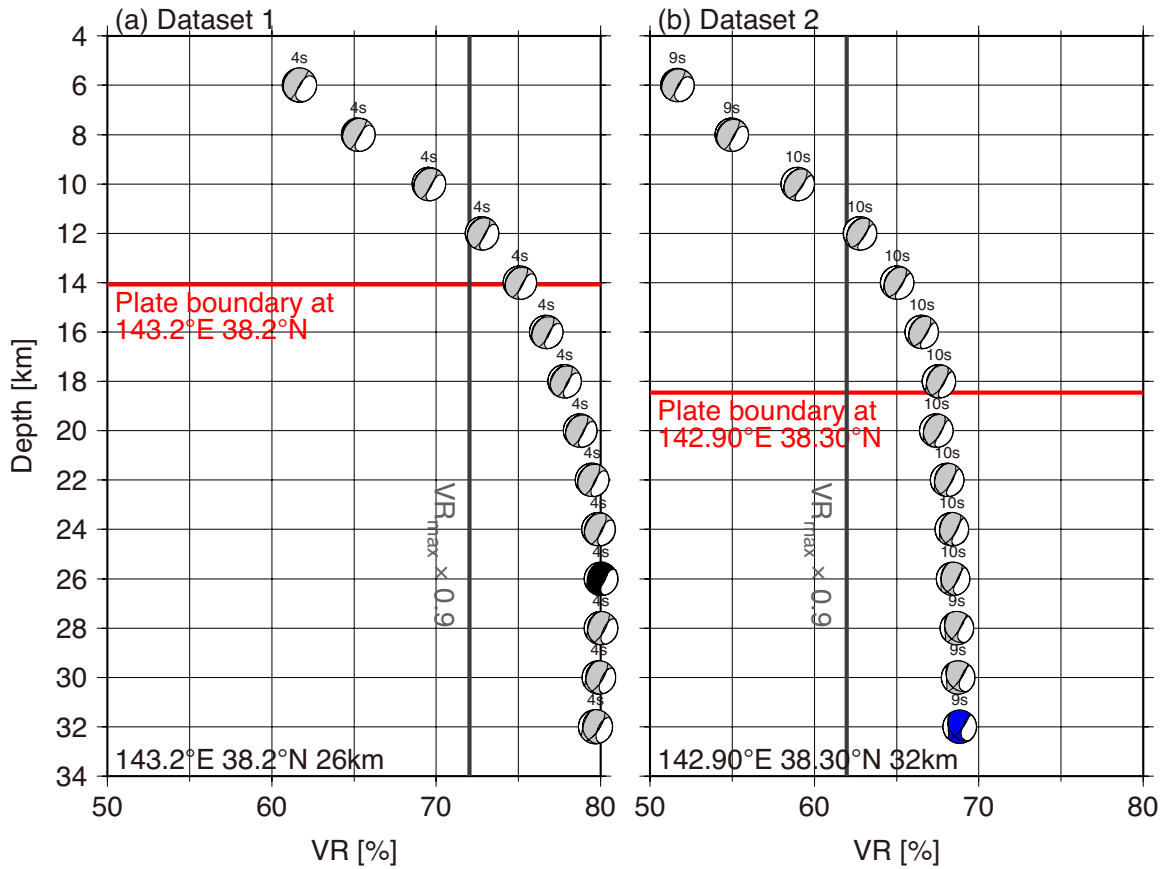


Figure S2. Vertical distribution of VR and moment tensors at the horizontal location of the best-fit CMT solution for foreshock #2 (Figure 5). (a) Result from dataset consisting of only onshore seismograms (Figure 5a). (b) Result from dataset consisting of onshore seismograms and offshore pressure data (Figures 5b). The location of the best-fit centroid is shown in the bottom left in each figure. Horizontal and vertical axes denote VR and centroid depth, respectively. Small numbers above each solution are the centroid time delay from the focal time. Red lines denote plate boundary depths obtained from seismic surveys by Ito et al. (2005).

Table S1. Locations of OBPGs

Station	Latitude [°N]	Longitude [°E]	Depth [m]
GJT3 ^a	38.2945	143.4814	3,293
P02 ^a	38.5002	142.5016	1,104
P03 ^a	38.1834	142.3998	1,052
P06 ^a	38.6340	142.5838	1,254
P07 ^a	38.0003	142.4488	1,059
P08 ^a	38.2855	142.8330	1,418
P09 ^a	38.2659	143.0006	1,556
TM1 ^b	39.2330	142.7830	1,564
TM2 ^b	39.2528	142.4500	954

^aPop-up recovery OBPG identical to those used in Hino et al. (2014) and Kubota et al. (2017)

^bReal-time cabled observation systems operated by the Earthquake Research Institute (ERI) of the University of Tokyo (Kanazawa & Hasegawa, 1997)

Table S2. Structure model used in this study^a

Depth [km]	Thickness [km]	P-wave velocity [km/s]	S-wave velocity [km/s]	Density [kg/m ³]	Qp	Qs
0	3	5.50	3.14	2300	600	300
3	15	6.00	3.55	2400	600	300
8	15	6.70	3.83	2800	600	300
18	67	7.80	4.46	3200	600	300
33	125	8.00	4.57	3300	600	300
100	100	8.40	4.80	3400	600	300
225	100	8.60	4.91	3500	600	300
425	—	9.30	5.31	3700	600	300

^aThis structure is same as that used in the F-net MT calculation and considered to be suitable for the one-dimensional structure of inland Japan (Kubo et al., 2002).